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## On the Theorem of Schur in the Special Cartan Space.

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矢野晋平：特殊 Cartan 空間に於ける Schur の定理について

The purpose of this paper is to explain that the theorem of *Schur* in the Riemannian geometry [1] holds good again for the special *Cartan* space in which the fundamental conditions are given, that is

$$(1) \quad \Gamma_{jk}^{*i} \parallel^l = 0, \quad R_{ijk} = 0.$$

L. BERWALD showed in his paper [2], that the necessary and sufficient conditions under which the Cartan space is the same as *Minkowski* space are that  $\Gamma_{jk}^{*i} \parallel^l = 0$  and  $R_{ijkl} = 0$ . Accordingly, our space may be said a more general space than the *Minkowski* space.

§ 1. *The curvature tensor of the space.* From the fundamental conditions (1) and the identities explained by L. BERWALD (Acta, (15. 1), (17. 1), (12. 12), (12. 7), (12. 10), (18. 8) and (18. 10)) we have

$$(2) \quad \bar{R}_{ihk}^j = 2\Gamma_{i[hk]}^{*j} + 2\Gamma_{i[h}^{*a} \Gamma_{|a|k]}^{*j},$$

$$(3) \quad R_{ihk}^j = \bar{R}_{ihk}^j, \quad P_{ijh}^k = 0,$$

$$(4) \quad A_{ij}^k \parallel_h = 0, \quad R_{ij(hk)l} = 0,$$

$$(5) \quad R_{ijhk} \parallel^l - 2R_{ij(a:h} A_k)^{al} = 0,$$

where  $R_{ijhk} \parallel^l = R_{ijhk} \parallel^l - R_{ajhk} A_i^{al} - R_{iahk} A_j^{al} - R_{ijak} A_h^{al} - R_{ijha} A_k^{al}$ ,

$\partial \Gamma_{ijh}^{*j} / \partial x^k = \Gamma_{ijk}^{*j}$ . Throughout this paper we shall use this notation: The first relation of (4) is rewritten in the following equation

$$(6) \quad A_{ij,h}^k + A_{ij}^k \parallel^a \Gamma_{aoh}^{*} - A_{aj}^k \Gamma_{ih}^{*a} - A_{ia}^k \Gamma_{jh}^{*a} + A_{ij}^a \Gamma_{ah}^{*k} = 0.$$

From (2), (3) and the second condition of (1), considering the relation (Acta, (8. 6))

$l_j, k = l_j \Gamma_{aok}^{*} (l^a - A^a)$ , we have

$$(7) \quad \Gamma_{io[k, l]}^* = \Gamma_{i[k}^{*a} \{l_a(l^\beta - A^\beta)\Gamma_{\beta o|l]}^* - \Gamma_{|ao|l]}^* \}.$$

Considering the known relation (Acta, (5.14))  $(L/\sqrt{g})^i = L/\sqrt{g}(l^i - A^i)$ , for any kind of tensor  $T_r^a \dots \beta_\delta$ , we have a general relation  $T_r^a \dots \beta_\delta^{[j k]} = T_r^a \dots \beta_\delta^{[j(l^k) - A^k]}$ . If we put the metric tensor  $g_{ij}$  in it, we have

$$(8) \quad A_{ij}^{[a \parallel \beta]} = A_{ij}^{[a l \beta]} - A^{\beta]}$$

If the obtained equation (Acta (6. 8))  $\Gamma_{i k}^{*j} = \Gamma_{i k}^j + A_i^j \Gamma_{a o k}$ , be multiplied by  $l_j$  and  $(\delta_h^i - l_h A^i)$ , and  $i$  and  $j$  be contracted, considering the relation  $(\delta_i^a + l_i A^a)(\delta_h^i - l_h A^i) = \delta_h^a$ , we may solve the  $\Gamma_{h o k}$  conversely, that is

$$(9) \quad \Gamma_{h o k} = \Gamma_{i o k}^{*i} (\delta_h^i - l_h A^i).$$

If we insert  $\Gamma_{h o k}$  into the known equations (Acta (7. 1), (7.3)) such as

$$g_{ij, k} = 2\Gamma_{(ij)k}^* - 2A_{ij}^a \Gamma_{a o k} \quad \text{and} \quad \Gamma_{ijk}^* = \gamma_{ijk} + A_{ij}^a \Gamma_{a o k} - A_{ik}^a \Gamma_{a o j} + A_{jk}^a \Gamma_{a o i}$$

where  $\gamma_{ijk} = \frac{1}{2}(g_{ij, k} + g_{jk, i} - g_{ik, j})$ ,

we obtain again the formal equations

$$(10) \quad g_{ij, k} = 2\Gamma_{(ij)k}^* - 2A_{ij}^a \Gamma_{a o k}^*$$

$$(11) \quad \Gamma_{ijk}^* = \gamma_{ijk} + A_{ij}^a \Gamma_{a o k}^* - A_{ik}^a \Gamma_{a o j}^* + A_{jk}^a \Gamma_{a o i}^*$$

If we put the equations (10), (11) in the relation (2), in the following form,

$$A_{ij}^a \Gamma_{a o k}^* - A_{ik}^a \Gamma_{a o j}^* + A_{jk}^a \Gamma_{a o i}^* = T_{ijk}, \quad \text{we have}$$

$$(12) \quad R_{ijkl} = 2(\gamma_{ij[h, k]} + T_{ij[h, k]} + 2A_{aj}^\beta \Gamma_{\beta o[k}^* \Gamma_{h]i}^{*a} - \Gamma_{i[h}^{*a} \Gamma_{k]aj}^*)$$

Considering the relation  $\gamma_{ij[h, k]} = \gamma_{hk[i, j]}$  we have the following equations

$$R_{ijhk} - R_{hkij} = 2(A_{ij, [k}^a \Gamma_{|ao|h]}^* + A_{ij}^a \Gamma_{ao[h, k]}^* - A_{i[h, k]}^a \Gamma_{a o j}^* - A_{i[h}^a \Gamma_{|a o j, k]}^* + A_{j[h, k]}^a \Gamma_{a o i}^* + A_{j[h}^a \Gamma_{|a o i, k]}^* + 2A_{aj}^\beta \Gamma_{\beta o[k}^* \Gamma_{i]h}^{*a}) \text{ cycle}(ih)(jk).$$

If we put the relations (6), (7) and (8) in the above equations and after lengthy calculations, we find that the right side is zero. Thus we have the

*Theorem.* In the special Cartan space given the fundamental conditions (1), we have the same relation as in the Riemannian space

$$(13) \quad R_{ijhk} = R_{hkij}.$$

§ 2 *The theorem of Schur.* As in the Riemannian geometry we define the curvature  $K$  at a point and an element of hypersurface  $u_i$  for the orientation determined by two vectors  $V_1^i$  and  $V_2^k$  as follows

$$K = R_{ijkl} V_1^i V_2^j V_1^k V_2^l / (g_{ik}g_{jl} - g_{il}g_{jk}) V_1^i V_2^j V_1^k V_2^l.$$

From (13) it follows that a necessary and sufficient condition under which the curvature at every point and element of hypersurface of our space be independent of the orientation is that

$$(14) \quad R_{hijl} = -K(g_{hk}g_{jl} - g_{hl}g_{jk}).$$

If we differentiate covariantly the equation (14) with  $i$  and put in the second relation of (4), and multiply by  $g^{hk}$  and  $g^{jl}$ , and contract the same indexes, we have the relation

$$(15) \quad K|_i = K_{,i} + k_{||}^a \Gamma_{aoi}^* = 0.$$

Similarly, differentiate with  $u_i$ , and put in (5), we have

$$K^{|l} (g_{ih}g_{jk} + g_{ih}g_{kh}) + 2R_{ija(h} A_k)al = 0.$$

If the above equation be multiplied by  $l^i$  and  $i$  be contracted, from the second condition of (1) we have

$$(16) \quad K^{|l} = (L/\sqrt{g}) \partial K / \partial u_l = 0.$$

From (15) and (16) we have finally  $K = const.$  Thus we have the

*Theorem. The theorem of Schur in the Riemannian geometry holds good again in the special Cartan space in which the fundamental conditions (1) are given.*

(October, 1951)

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## On the Molecular Weight of Liquid Polymers.

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清水 清 : 液体重合体の分子量に就て

§ 1 Sound velocity in liquid became useful for the investigation of molecular structure after the discovery by M. R. Rao<sup>1)</sup> of the empirical constant (called the molar sound velocity)

$$R = M \frac{v^{\frac{1}{3}}}{d} \quad (1)$$

where  $M$ ,  $v$ , and  $d$  are the molecular weight, sound velocity, and density.

Further, R.T. Lagemann and W.S. Dunbar<sup>2)</sup>

showed that, within a homologous series of liquid polymers, linear relationships exist between any two molar constants. An example of such a relationship is the one involving molar refraction  $N$  and molar sound velocity  $R$

$$R = AN + B \quad (2)$$

where  $A$  and  $B$  are, respectively, the slope and intercept.

Recently, A. Weissler, S. W. Fitzgerald,