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## on a special Kawaguchi space

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北村五郎；特殊河口空間について

The fundamental theory of  $n$ -dimensional spaces such that the arc length of a curve  $x^i = x^i(t)$  is given by the integral  $s = \int \left\{ A^i(x, x') x'^{i1} + B(x, x') \right\}^{\frac{1}{p}} dt$  has been established by Prof. Kawaguchi<sup>1)</sup> and thereafter studied by many others.

In the present paper we study a generalization of an absolute derivative with  $p \neq 1, 3$  when the components of a vector  $v^i$  are homogeneous of degree  $h$  with regard to the  $x^i$  (§ I) and lead a necessary and sufficient condition that the special Kawaguchi space of even dimensions with  $|H_{ij}| \neq 0$  be locally flat by another method from Watanabe's proof<sup>2)</sup> (§ 2).

The notations are the same as those of Kawaguchi's paper.

## § I. A generalization of an absolute derivative.

The definitions of the absolute derivative

$$\delta v^i = dv^i + \Gamma_{(j)(k)}^i v^j dx^k$$

have no geometrical meaning if  $v^i$  are homogeneous of degree  $h$  with regard to the  $x$ . For we obtain

$$(\delta v^i)_{\bar{t}} = \sigma^h \delta v^i + h v^i \sigma^{h-1} d\sigma \quad \left( \sigma = \frac{dt}{d\bar{t}} \right)$$

by transformation of the parameter  $t : t = t(\bar{t})$ .

When we put

$$\delta^* v^i = \delta v^i - h v^i f(t),$$

we get

$$(\delta^* v^i)_{\bar{t}} = \sigma^h \{ \delta v^i - h v^i (f)_{\bar{t}} \} + h v^i \sigma^{h-1} d\sigma$$

by transformation of  $t$ .

Therefore  $f$  must have the following form by transformation  $t$

$$(f)_{\bar{t}} = f + \frac{d\sigma}{\sigma}.$$

On the other hand

$$(\delta A_i x^{(2)i})_{\bar{t}} = \sigma^p \delta A_i x^{(2)i} + (p-2) F \sigma^{p-1} d\sigma - \sigma^{p-2} \frac{d\sigma}{d\bar{t}} A_i \delta x'^i,$$

$$(T_i \delta x'^i)_{\bar{t}} = \sigma^p T_i \delta x'^i + p F \sigma^{p-1} d\sigma - (2p-3) \sigma^{p-2} \frac{d\sigma}{d\bar{t}} A_i \delta x'^i.$$

Accordingly we get

$$\left( \frac{\delta A_i x^{(2)i}}{p F} \right)_{\bar{t}} = \frac{\delta A_i x^{(2)i}}{p F} + \frac{p-2}{p} \frac{d\sigma}{\sigma} - \frac{1}{p F} \sigma^{-2} \frac{d\sigma}{d\bar{t}} A_i \delta x'^i,$$

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$$\begin{aligned} \left( \frac{T_i \delta x'^i}{(2p-3)pF} \right)_{\bar{t}} &= \frac{T_i \delta x'^i}{(2p-3)pF} + \frac{1}{2p-3} \frac{d\sigma}{\sigma} - \frac{1}{pF} \sigma^{-2} \frac{d\sigma}{d\bar{t}} A_i \delta x'^i, \\ \therefore \left( \frac{(2p-3) \delta A_i x'^{2i} - T_i \delta x'^i}{(2p-3)pF} \right)_{\bar{t}} &= \frac{(2p-3) \delta A_i x'^{2i} - T_i \delta x'^i}{(2p-3)pF} \\ &\quad + \left( \frac{p-2}{p} - \frac{1}{2p-3} \right) \frac{d\sigma}{\sigma}. \end{aligned}$$

Consequently we can put

$$f = \frac{(2p-3) \delta A_i x'^{2i} - T_i \delta x'^i}{2(p^2+1)F}.$$

Using the relations :  $x'^{2;k} = -T_i G^{ik}$ ,  $G_{ik} = 2A_{i(k)} - A_{k(i)}$ ,  
we put

$$\begin{aligned} \lambda_i &= \frac{2p-3}{2(p^2+1)F} x'^{2;k} \nabla_i A_k, \\ \mu_i &= \frac{1}{(p^2+1)F} x'^{2;k} \left\{ (p-2)A_{k(i)} + A_{i(k)} \right\}. \end{aligned}$$

Then we obtain as the required absolute differential

$$\delta^* v^i = \delta v^i - h v^i (\lambda_j dx^j + \mu_j \delta x'^j),$$

accompanied by the covariant derivatives

$$\begin{aligned} \nabla_j^* v^i &= \nabla_j v^i - \frac{2p-3}{2(p^2+1)F} h v^i x'^{2;k} \nabla_j A_k, \\ \nabla_{j'}^* v^i &= \nabla_{j'} v^i - \frac{1}{(p^2+1)F} h v^i x'^{2;k} \left\{ (p-2)A_{k(j)} + A_{j(k)} \right\}. \end{aligned}$$

$F^{-\frac{h}{p}} \delta^* v^i$  are independent of a change of the parameter  $t$ .

§ 2. Another proof of the theorem.

Theorem : A necessary and sufficient condition that the special Kawaguchi space of even dimensions with  $|H_{ij}| \neq 0$  be locally flat is that the following relations are satisfied :

$$A_{i(j)(k)} = 0, \quad R_{jki}^{\dots i} = 0, \quad B_{jki}^{\dots i} = 0, \quad \delta A_{i(j)} = 0.$$

Proof. Necessity.

If the space be locally flat, the relations  $A_{a(b)} = \text{const.}$ ,  $B(y, y') = 0$  are satisfied in the suitable chosen  $y$  coordinate system.

Hereafter we use the indexes  $i, j, k, \dots$ ;  $a, b, c, \dots$  in regard to  $x$  and  $y$  coordinate system respectively.

We can lead easily the relation  $A_{i(j)(k)} = 0$  from  $A_{a(b)} = \text{const.}$ .

Differentiating  $A_{a(b)} y'^b = (p-2)A_a$  with respect to  $y'^c$ , we get  $p=3$ .

Consequently we have the following relations :

$$A_{a(b)} y'^b = A_a, \quad A_{b(a)} y'^b = -A_a, \quad \therefore A_{b(a)} = -A_{a(b)}.$$

Besides

$$\begin{aligned} H_{ab} &= 2A_{a(b)}, \quad G_{ab} = 3A_{a(b)} = \text{const.}, \\ 2\Gamma^c &= (2A_{ab} y'^b - B_{(a)}) G^{ac} \\ &= (A_a G^{ab})_b y'^b = 2 \left( -\frac{y'^c}{2p-3} \right)_b y'^b \\ &= 0. \end{aligned}$$

Hence we have  $R_{jki}^{\dots i} = 0$ ,  $B_{jki}^{\dots i} = 0$  from  $\Gamma^c = 0$  and  $\delta A_{a(b)} = 0$  from  $A_{a(b)} = \text{const.}$

Sufficiency.

From  $B_{jki}^{\dots i} = 0$ , then  $\Gamma_{(j)(k)}^i$  are functions only of the  $x'$  s. If we consider now the simultaneous partial differential equations

$$\begin{cases} \frac{\partial y^a}{\partial x^i} = P_i^a \\ \frac{\partial P_i^a}{\partial x^j} = P_m^a \Gamma_{(i)(j)}^m, \end{cases}$$

these equations satisfy the integrability conditions from  $R_{jki}^{\dots i} = 0$ ,  $B_{jki}^{\dots i} = 0$ .

Hence we choose the above solution  $y^a = y^a(x)$  as the coordinate transformation.

Differentiating  $P_i^a Q_a^j = \delta_i^j \left( Q_a^i = -\frac{\partial x^j}{\partial y^a} \right)$  by  $y^a$ , we have

$$Q_{aa}^j = -Q_a^i Q_b^j Q_m^a P_{im}^a = -Q_a^i Q_m^a \Gamma_{(i)(m)}^j.$$

Differentiating  $A_{a(b)} = Q_a^i Q_b^j A_{i(j)}$  by  $y^c$  and using the above relations, we can see

$$\begin{aligned} A_{a(b)c} &= (Q_{ac}^i Q_b^j + Q_a^i Q_{bc}^j) A_{i(j)} + Q_a^i Q_b^j Q_c^k A_{i(j)k} \\ &= Q_a^i Q_b^j Q_c^k \nabla_k A_{i(j)} \\ &= 0, \\ B(y, y') &= B(x, x') - A_m Q_a^m P_{ij}^a x'^i x'^j \\ &= B(x, x') - A_m \Gamma_{(i)(j)}^m x'^i x'^j \\ &= 0. \end{aligned}$$

Hence this space be locally flat.

### References

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